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COMMENT

**False eigenvalues of the Hill determinant method**

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**Abstract.** It is shown analytically that the false eigenenergies of a potential obtained by the method of Hill determinants are the physical energies of a different potential.

Recently, Killingbeck (1986) has observed numerically a relationship between the false eigenvalues, as obtained by the Ginsberg-Killingbeck (Ginsberg 1982, Killingbeck 1986) version of the Hill determinant method, for the potential  $-r^{-1} + r + r^2$  and the true energies for the potential  $r^{-1} - r + r^2$ . In the present comment we have demonstrated analytically that the following is a general rule. Given a potential  $V(r) = a/r + br + cr^2$  when  $c > 0$ , those eigenvalues identified as false by Killingbeck's method are in fact true eigenvalues of the associated potential  $V''(r) = -a/r - br + cr^2$ .

For the perturbed Coulomb potential

$$V(r) = a/r + br + cr^2 \quad c > 0. \tag{1}$$

Killingbeck has shown numerically that the false eigenvalues reveal themselves by having negative values for the expectation values of  $r$  and  $r^{-1}$  obtained by the technique developed by him (Killingbeck 1985). If a potential term  $\lambda r^n$  is added to the Hamiltonian the change in energy according to the first-order perturbation theory is

$$\delta E = \lambda \langle r^n \rangle \tag{2}$$

where  $\lambda$  is quite small. Killingbeck (1986) has applied this equation to compute  $\langle r \rangle$  and  $\langle r^{-1} \rangle$  for the potential (1) and obtained the negative results for some values of the parameters. (For a more detailed explanation of the  $\langle r^n \rangle$  computations the reader is referred to equation (5) of Killingbeck (1986).) We explain here why the Hill determinant method yields such unphysical results.

By applying the simple power-series ansatz

$$\psi(r) = r^{l+1} \exp(-\beta r^2 - \alpha r) \sum_{n=0}^{\infty} A_n r^n \tag{3}$$

to the radial Schrödinger equation

$$-D^2\psi + [V(r) + l(l+1)/r^2]\psi = E\psi \tag{4}$$

the differential equation (4) is transformed to the difference equation

$$b_n A_{n+1} + a_n A_n + c_n A_{n-1} + d_n A_{n-2} = 0 \quad n \geq 0 \tag{5}$$

with

$$\begin{aligned} b_n &= (n+1)(n+2l+2) \\ a_n &= -2\alpha(n+l+1) - a \\ c_n &= -2\beta(2n+2l+1) + \alpha^2 + E \\ d_n &= 4\alpha\beta - b \\ \beta &= \frac{1}{2}\sqrt{c} > 0 \quad A_{-1} = A_{-2} = 0. \end{aligned}$$

The eigenvalue condition of the Hill determinant for large  $N$  is

$$\det Q_N = 0 \quad (6)$$

$$Q_N = \begin{pmatrix} a_0 & b_0 & 0 & 0 & \dots \\ c_1 & a_1 & b_1 & 0 & \dots \\ d_2 & c_2 & a_2 & b_2 & \dots \\ \vdots & & & & \\ 0 & \dots & 0 & d_{N-1} & c_{N-1} & a_{N-1} \end{pmatrix}. \quad (7)$$

It may be noted that

$$A_N = (-1)^N A_0 \det Q_N / (b_0 b_1 \dots b_{N-1}) \quad (8)$$

which shows that the vanishing of  $\det Q_N$  ensures the vanishing of  $A_N$ . Killingbeck (1986) has computed the eigenvalues from the zeros of  $A_N$  in terms of the parameter  $E$  for some large  $N$  and obtained the physical eigenenergies for the potential (1) with a sufficiently large positive  $\alpha$ . As soon as the sign of  $\alpha$  is reversed and made negative the unphysical eigenvalues are produced in this method. When we change  $\alpha \rightarrow -\alpha$  in equation (7) we get the determinant  $Q'_N$  which is completely different from  $\det Q_N$  and therefore we get different eigenvalues from the zeros of the two determinants. However, in (7) if we do not change the sign of  $\alpha$ , but change the signs of  $a$  and  $b$ , i.e.  $\alpha \rightarrow \alpha$ ,  $a \rightarrow -a$  and  $b \rightarrow -b$  we get the determinant  $Q''_N$  with the following property:

$$\det Q''_N = (-1)^N \det Q'_N. \quad (9)$$

Thus the false eigenvalues of the potential  $a/r + br + cr^2$  obtained from  $\det Q'_N = 0$  are the physical eigenvalues of the potential  $-a/r - br + cr^2$  obtained from  $\det Q''_N = 0$ . Killingbeck (1986) discovered this finding by the process of computation. Here we establish the result analytically.

Killingbeck (1986) has applied equation (2) for computation of  $\langle r^n \rangle$  from  $\det Q'_N = 0$  which gives negative values of  $\langle r \rangle$  and  $\langle r^{-1} \rangle$ . It may be noted that the zeros of  $\det Q''_N$  are identically the same as those of  $\det Q'_N$  leading to the same value of  $\delta E$ , but the sign of  $\lambda$  in equation (2) is reversed. Thus the false eigenvalues of the table 1 of the paper of Killingbeck (1986) should be reinterpreted accordingly. If  $x$  and  $y$  are the expectation values of  $r$  and  $r^{-1}$  as obtained from (2) for the potential  $-r^{-1} + r + r^2$  in case of false eigenvalues derived from the equation  $\det Q'_N = 0$ , the false energies agree with the true energies for the potential  $r^{-1} - r + r^2$  and the expectation values of  $r$  and  $r^{-1}$  for the potential  $r^{-1} - r + r^2$  according to equation (2) are  $-x$  and  $-y$  giving positive sign to the expectation values when  $x$  and  $y$  are negative.

## References

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